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Research Article

Studying the nuclear structure for the (70Se, 70Ga, 70Zn) isotopes using NuShellX@MSU code

Sara F. Hadi^{1*}, Ali K. Hasan²

¹University of Kufa, Faculty of Education for Girls, Physics Department, Najaf, Iraq. * Corresponding Author Email: sarahf.nehma@student.uokufa.edu.iq ORCID: 0009-0004-9069-6693

²University of Alkafeel, Najaf, Iraq.

Email: alikh.alisnayyid@alkafeel.edu.iq ORCID: 0000-0002-8126-5179

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Abstract:

The structure of neutron-rich even-even isotopes of A collection of isotopes possessing an identical mass number (A=70), specifically (70Se, 70Ga, 70Zn) has been examined by extensive shell-model simulations. The probabilities B(E2) and B(M1) are calculated using the shell model with NushellX@MSU code, utilizing the effective interaction f5pvh. The findings for excitation energies and reduced transition Probabilities are juxtaposed with the most recent experimental data available. A satisfactory concordance is achieved for all examined isotopes.

1. Introduction

The nuclear shell model was proposed almost 70 years ago by Maria Goeppert Mayer and Hans Jensen, who elucidated the stability of specific nuclei via the notion of "magic numbers" [1]. These investigations have questioned the universality of conventional magic numbers, demonstrating that shell structure can fluctuate considerably with variations in proton and neutron counts[2].Initial implementations of the shell model utilised empirical effective Hamiltonians to characterise nuclear structure, concentrating on nuclei close to stability[1]. Methods such as the Lanczos algorithm and the Monte Carlo shell model have been devised to approach accurate diagonalisation, facilitating the examination of medium- and heavy-mass nuclei[3]. The nucleus constitutes a quantum mechanical configuration, consisting of protons neutrons[4][5]. The behaviour of these nuclei is affected by the movement of nucleons situated outside the dense central core.made of valence particles[6]. A precise comprehension of the interactions of valence particles is essential for effectively elucidating the characteristics of the nucleus. A variety of nuclear models have clarified the properties of several nuclei[7][8]. Among these models, the distinguished nuclear shell model is prominent. This model delineates the energy states,

their configurations, and the transitions occurring them, so offering a representation of nuclear characteristics[9,10]. The nuclear shell model has proven highly useful in elucidating nuclear structure; upon identifying an appropriate effective interaction, the shell model routinely and reliably predict may observables[11]. Principal elements: the interaction among nucleons (N-N interaction) and the spatial configuration designated for valence particles. One can do shell-model calculations utilising either a true N–N interaction within a broad configuration space or a modified effective interaction within a more limited configuration space [12]. Owing to shortrange correlations and medium effects, genuine nucleon-nucleon (NN) interactions necessitate renormalisation for application in shell-model computations[13]. This type of engagement is referred to as a "active interaction" [14]. The KB3G [15] and GXPF1A Hamiltonians have yielded predictions for the spectra in this domain, serving as a benchmark for several investigations over the past two decades[16][17]. Both of these represent "universal" Hamiltonians for the p-f model space. Recent findings indicate that a data-driven Hamiltonian for calcium isotopes enhances the representation of all existing data[18]. For light nuclei, there exist multiple "standard" effective interactions, including Cohen-Kurath the

interaction[19]. The isotopes possessing an identical mass number (70Se; Z = 34, 70Ge; Z=32, 70Zn; A=30) encompass three doubly-closed shells with neutron numbers ranging from N = 36 to N = 40, which have been analyzed through advanced shell model calculations utilizing recently developed interaction. Employing(56Ni) as a core within the (f5/2-p p1/2) model space provides a distinctive framework for examining the evolution of shell structure. Isotopic isotopes have garnered new scientific interest to address the question of magicity vs superfluidity about the doubly magic nature of these nuclei[20-24]. Srivastava [25] conducted model computations of the shell for zinc and other isotopes by altering the fpg reaction, adjusting 28 elements of the two-body matrix from the prior reaction. The revised reaction, termed fpg9a, was calibrated for copper isotopes and evaluated for zinc and nickel isotopes. F. Recchia et al. [26] recently examined the level structure of 68Ni using twoneutron ejection and a multiple nucleon transfer process. They juxtapose their experimental findings with shell model calculations employing various contemporary effective interactions. Y. Tsunoda[27] investigated the configurations of neutron-rich exotic many of isotopes using extensive shell model computations utilizing the improved Monte Carlo shell model (MCSM), in which the experimental energy levels are accurately represented using a singular fixed Hamiltonian. This research presents shell model calculations in the f5p-shell region for the even-even (70Se, 70Ga, 70Zn) isotopes, utilizing the contemporary f5pvh effective interaction to evaluate their capacity to replicate experimental results in this mass range [28,29].

2. Shell model

The independent-particle Hamiltonian of an Aparticle system can be expressed in terms of twoparticle interactions as [16]:

$$= \sum_{k=1}^{A} T_k + \sum_{\substack{k=1\\k=1}}^{A} \sum_{l=k+1}^{A} W(\vec{r}_k, \vec{r}_l)$$
 (1)

where W(rk, rl) \) represents the two-body interaction between the kth and lth nucleons. By selecting an average potential U(r_k), the Hamiltonian is expressed as [16].

$$.H = \sum_{k=1}^{A} [T_k + U(r_k)] + \sum_{k=1}^{A} \sum_{l=k+1}^{A} W(\overrightarrow{r_k}, \overrightarrow{r_l}) - \sum_{k=1}^{A} U(r_k)$$
 (2)

The initial term corresponds to the independentparticle Hamiltonian, whereas the subsequent second and third terms address the divergence from independent particle motion, referred to as the residual interaction. By partitioning the summations into core and valence components, equation (2) can be reformulated [20].

$$H = H_{core} + H_1 + H_2 + V(\overrightarrow{r_1}, \overrightarrow{r_2})$$
 (3)

In the aforementioned equation, H_{core} encompasses all interactions among the nucleons constituting the core, H₁ and H₂ represent the single-particle contributions from particles 1 and 2, respectively, while $V(\vec{r_1},\vec{r_2})$ denotes the residual interaction that characterizes all interactions between particles 1 and 2, as well as any interactions with core nucleons. Substituting this variant of the Hamiltonian into the Schrödinger equation produces a corresponding expression for the energy [16].

$$E = E_{core} + E_1 + E_2 + \langle \Phi_{J,\tau} | V(\vec{r_1}, \vec{r_2}) | \Phi_{J,\tau} \rangle$$
(4)

Here, E_{core} represents the binding energy of the core nucleus, while E_1 and E_2 denote the single-particle energies of the orbitals external to the core. Additionally, $\langle \Phi_{I,\tau} | V(\overrightarrow{r_1}, \overrightarrow{r_2}) | \Phi_{J,\tau} \rangle$ signifies the residual interaction that must be theoretically specified. It is essential to recognize that the energy provided by equation (4) pertains exclusively to pure arrangements. In principle, every adjacent state possessing identical total angular momentum J and total isospin τ will undergo mixing. The mixed eigenstates are represented by linear combinations of the unperturbed wave functions [1,16].

$$\left(\psi_{J,\tau}\right)_p = \sum_{k=1}^g a_{kp}(\Phi_{J,\tau})_p \tag{5}$$

where g represents the number of mixed configurations, and the label p takes values from 1 to g. The coefficients a_{kp} satisfy the condition [4].

$$\sum_{k=1}^{g} \left| a_{kp} \right|^2 = 1 \tag{6}$$

Substituting equation (5) into the Schrödinger equation yields,

$$H(\psi_{J,\tau})_p = E_p(\psi_{J,\tau})_p$$
 (7) resulting in a system of linear equations [20].

3. Results and Discussions

3.1Excitation Energies

The core is represented by (56Ni) for all studied isotopes, with valence nucleons allocated in the (1f5/2, 2p3/2, and 2p1/2) valence space, utilizing the f5pvh effective interaction through the shell model NushellX@MSU code. The computed energy levels for (70Se, 70Ga, 70Zn) isotopes are compared utilizing f5pvh effective interaction alongside the most recent experimental data. The theoretical calculations encompass fourteen nucleons external to the closed core of the (70Se, 70Ga, 70Zn) nucleus, consistent with the anticipated conclusions presented in The tables (1,2 and 3) regarding total angular momentum and symmetry

Table (1) delineates a comparison between the experimental outcomes for the 70Se isotope utilizing the f5pvh interaction and the presently accessible theoretical conclusions (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Data Sheets 136, 1 (2016) 10.1016/j.nds.2016.08.001" n.d.).

1. 70Se

Table 1. Excitation energy predictions for the 70Se isotope using f5pvh interaction and observed experimental energies comparison

	cal values of (MeV)	Experimental values		
J ⁺	f5pvh	E(MeV)	J	
01	0	0.0	0+	
21	0.873	0.944	2+	
02	1.534			
22	1.591	1.599	2+	
23	2.019	2.010	(0+)	
41	2.083	2.038	4+	
31	2.196			
24	2.245			
25	2.445			
42	2.473	2.382	4+	
11	2.566	2.553		
32	2.586	2.518	4(-)	
03	2.593			
43	2.691			
26	2.711			
04	2.713			
12	2.766			
44	2.876			
27	2.895			
45	2.917			
33	2.928			
34	2.984			
46	3.051			
13	3.074			
35	3.126	3.139		
47	3.163			
28	3.186			
29	3.205			
36	3.289			
51	3.299	3.218	(6 ⁺)	
210	3.315			
14	3.316			
52	3.335	3.356		
05	3.417			
15	3.418			
37	3.465			
48	3.518	3.524	(5-)	

49	3.566		
06	3.569		
61	3.621	3.788	(6-)
	3.626	3.766	
410			
16	3.645	3.644	
38	3.679		
17	3.733		
39	3.802		
18	3.813		
53	3.883		
310	3.887		
62	3.893		
19	3.937		
07	3.949		
54	4.033		
110	4.071		
63	4.112		
64	4.162		
08	4.191		
0_{9}	4.217		
55	4.234		
56	4.261		
010	4.328	4.324	
57		4.324	
	4.347		
58	4.374		
65	4.398	4.410	
66	4.58		
59	4.628		
510	4.663		
67	4.709		
68	4.785		
71	4.887	4.896	(9-)
69	5.131		
72		5.209	(9-)
	5.21	5.209	(9-)
610	5.21 5.233		
6 ₁₀ 7 ₃	5.21 5.233 5.373		
6 ₁₀ 7 ₃ 7 ₄	5.21 5.233 5.373 5.44		
6 ₁₀ 7 ₃ 7 ₄ 8 ₁	5.21 5.233 5.373 5.44 5.532		
6 ₁₀ 7 ₃ 7 ₄ 8 ₁ 8 ₂	5.21 5.233 5.373 5.44 5.532 5.571		
6 ₁₀ 73 74 81 82 75	5.21 5.233 5.373 5.44 5.532		
6 ₁₀ 73 74 81 82 75	5.21 5.233 5.373 5.44 5.532 5.571		
6 ₁₀ 7 ₃ 7 ₄ 8 ₁ 8 ₂ 7 ₅ 7 ₆	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844		(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905	5.693	
6 ₁₀ 73 74 8 ₁ 8 ₂ 75 76 77 8 ₃	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01	5.693 	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039	5.693 6.017	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094	5.693 	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155	5.693 6.017	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094	5.693	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174	5.693	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177	5.693	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362	5.693	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467	5.693	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56	5.693 6.017 6.510	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68	5.693 6.017 6.510 6.602	(10 ⁺) (10 ⁺) (10 ⁺) (110 ⁺) (110 ⁺) (110 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812	5.693 5.693 6.017 6.510 6.602	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68	5.693 6.017 6.510 6.602	(10 ⁺) (10 ⁺) (10 ⁺) (110 ⁺) (110 ⁺) (110 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812	5.693 5.693 6.017 6.510 6.602	(10 ⁺)
6 ₁₀ 73 74 8 ₁ 8 ₂ 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079	5.693 5.693 6.017 6.510 6.602 6.956	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203	5.693 6.017 6.510 6.602 6.956	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331	5.693 5.693 6.017 6.510 6.602 6.956	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421	5.693 5.693 6.017 6.510 6.602 6.956 	(10 ⁺) (112 ⁺) (12 ⁺) (13 ⁻)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551	5.693 5.693 6.017 6.510 6.602 6.956 7.305	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421	5.693 5.693 6.017 6.510 6.602 6.956 	(10 ⁺) (112 ⁺) (12 ⁺) (13 ⁻)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551	5.693 5.693 6.017 6.510 6.602 6.956 7.305	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751	5.693 5.693 6.017 6.510 6.602 7.305	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97 102	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751 7.904	5.693 5.693 6.017 6.510 6.602 7.305 7.554	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97 102 98	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751 7.904 7.912	5.693 5.693 6.017 6.510 6.602 7.305 7.554 7.940	(10 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97 102 98 103	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751 7.904 7.912 7.986	5.693 5.693 6.017 6.510 6.602 7.305 7.554	(10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97 102 98	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751 7.904 7.912 7.986 8.025	5.693 5.693 6.017 6.510 6.602 7.305 7.554 7.940	(10 ⁺) (10 ⁺) (10 ⁺) (10 ⁺) (110 ⁺)
6 ₁₀ 73 74 81 82 75 76 77 83 78 79 84 85 710 86 87 88 89 810 91 92 93 94 101 95 96 97 102 98 103	5.21 5.233 5.373 5.44 5.532 5.571 5.682 5.844 5.905 6.01 6.039 6.094 6.155 6.174 6.177 6.362 6.467 6.56 6.68 6.812 6.977 7.079 7.203 7.331 7.421 7.551 7.688 7.751 7.904 7.912 7.986	5.693 5.693 6.017 6.510 6.602 7.305 7.554	(10 ⁺) (110 ⁺)

10_{4}	8.141		
105	8.19		
106	8.391	8.349	
107	8.67		
108	8.868		
109	8.972		
1010	9.072		
111	9.361		
112	9.386		
113	9.494	4.496	(16^+)
121	9.677		
114	10.037		
115	10.204		
116	10.354		
117	10.425		
118	10.557		
122	10.643	10.646	(18+)
119	10.717		
1110	10.759		
123	10.875		
124	11.061		
125	11.761		
126	12.094		
127	12.145		
128	12.323		
129	12.491		
131	12.502		
1210	12.615		
132	12.965		
133	14.379		
- 5		l	

By comparing the experimental data for this isotope in the table above with our theoretical findings using the f5pvh interaction, the following can be observed: A comparison with the practical numbers available indicated that the ground state parity and total angular momentum of the 0+ level were equal.

The angular momentum and parity of practical energy was determined to be (2.553, 3.139, 3.356, 3.644, 4.324, 4.410, 6.017, 8.349, 2.010, 3.218, 5.693, 6.510, 6.602, 6.956, 7.940, 4.496, 10.646) MeV, corresponding to positive parity angular momentum value of 1, 5, 1, 0, 6, 8, 2, 6, 10, 12, 12, 12,14, 16 and 18, respectively. And (3.524, 3.788, 4.896, 5.209, 8.017, 2.518) MeV with angular momentum 4, 6, 7, 7, 9, 9, 9, 3 but negative parity. This reflects the degree of alignment between the practical value and our theoretical value.

By comparing the theoretically computed energies (0.873 MeV; 2^+),) 1.591 MeV; 2^+),) 2.083; 4^+) and) 2.473 MeV; 4^+) with the available experimental data, we were able to get good agreement for the angular momentums.

According to our calculations, the maximum experimental energy value is 20.246MeV, and the greater predicted energy is theoretically 14.379MeV.

Through the theoretical calculations, we have (103) state with the total angular momentum & parity that have not symmetry by another practical value thus far.

2. 70Ga

Table (3) delineates a comparison between the experimental outcomes (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Data Sheets 136, 1 (2016) 10.1016/j.nds.2016.08.001" n.d.)for the 70Ga isotope utilizing the f5pvh interaction and the presently accessible theoretical conclusions .

Table 3. Excitation energy predictions for the 70Ga isotope using f5pvh interaction and observed experimental energies comparison (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Data Sheets 136, 1 (2016) 10.1016/j.nds.2016.08.001" n.d.)

Theoret	tical values	Experimental values			
J ⁺	E (MeV)	E (MeV)	J^{π}		
11	0.000	0.0	1+		
21	0.153				
41	0.276				
31	0.404				
32	0.447				
22	0.465	0.508	2+		
12	0.594	0.651	1+, 2+		
23	0.613	0.690	2-		
42	0.679	0.879	4-		
13	0.697				
01	0.738				
33	0.799				
14	0.818	0.901	1+, 2+, 3+		
24	0.888	0.995	2+		
02	0.989	1.002			
25	1.031	1.014	1+, 2+, 3+		
34	1.199	1.023	2+, 3+		
51	1.226	1.180	5		
35	1.279	1.253	3-, 4-		
26	1.281	1.258	1+ to 4+		
15	1.349	1.312	1+, 2+		
27	1.368	1.359	2+		
36	1.385				
28	1.400	1.445	1+, 2+		
03	1.420				

43	1.431		
16	1.470	1.456	1+, 2+
44	1.543	1.523	
37	1.587		
29	1.596	1.598	
17	1.630	1.633	1, 2, 3
210	1.682	1.691	
52	1.722	1.725	
38	1.736	1.735	
61	1.767		
45	1.812	1.807	
53	1.813	1.824	
18	1.880	1.846	
46	1.882	1.865	
04	1.888	1.877	
39	1.911	1.905	+
19	1.965	1.937	+
310	1.979	1.970	+
47	2.043	2.026	
110	2.151	2.142	
54	2.190	2.189	
48	2.301	2.300	+
49	2.331	2.351	
62	2.462	2.464	
55	2.495	2.477	
4 ₁₀	2.525	2.520	-
71	2.638	2.601	(8)
56	2.709	2.698	
05	2.748	2.726	
57	2.821	2.886	(9)
06	2.939		
63	2.954		
58	2.984		
59	3.089		
64	3.096		
510	3.221		
65	3.275		

07	3.414	
66	3.733	
72	3.739	
08	3.767	
09	3.809	
010	4.111	
67	4.193	
73	4.265	
81	4.292	
68	4.308	
69	4.489	
74	4.706	
610	4.743	
75	4.833	
76	6.347	
77	6.593	

By comparing the experimental data for this isotope in the table above with our theoretical findings using the f5pvh interaction, the following can be observed:

- A comparison with the available empirical data revealed that the ground state parity and total angular momentum of the 1+ level were identical.
- The angular momentum and parity of practical energies were determined to be $(1.002 \text{MeV}, 0^+)$ 1.523Mev, 4^{+} ,(1.598Mev, 2^+ , $(1.691 \text{MeV}, 3^+)$, (1.725Mev, 5^+),(1.735Mev,3^+), (1.807Mev, 4^{+}), $(1.824 \text{MeV}, 5^+), (1.846 \text{MeV}, 1^+), (1.865 \text{MeV},$ 4^+), $(1.877 \text{MeV}, 0^+)$, $(2.026 \text{MeV}, 4^+)$ $(2.142 \text{MeV}, [1] ^+), (2.189 \text{MeV}, 5^+), (2.351)$ 4^+), (2.464Mev, 6^+) ,(2.477Mev, Mev, 5^+),(2.726Mev, 5^+)and(2.698Mev,0 $^+$), corresponding to positive angular momentum. This reflects the degree of alignment between the practical value.
- By juxtaposing the theoretically calculated energies (0.465Mev, 2^+), (0.594Mev,1^+), (0.818Mev, [1] ^+), (0.888Mev, 2^+), (1.031Mev, 2^+), (1.199Mev,3^+), (1.281Mev, 2^+), (1.349Mev, 1^+), (1.400Mev, 2^+) and (1.470Mev, 1^+) with the existing experimental data, we achieved substantial concordance for the angular momentums and parity.
- This analysis anticipated that the total angular momentums for the practical energies of (1.905,

- 1.937, 1970, 2.300)MeV are consistent with parity matching.
- According to our calculations, the maximum experimental energy value is 8.349MeV, and the greater predicted energy is theoretically 14.379MeV.
- The total angular momentum was verified, and the parity was anticipated to be positive rather than negative at the experimental energies of 0.690, 0.879, and 1.253 MeV, corresponding to angular momenta of 2+, 4+, and 3+, respectively.
- We anticipate a total angular momentum for the indeterminate practical energies (2.601,2.886) MeV, corresponding to angular momenta of 7, 5 rather than 8,9 and a positive parity was determined for them.
- By juxtaposing the theoretically calculated energies with the existing experimental data, we achieved substantial concordance for the angular momentums for (1.633 MeV;1) and and symmetry of practical energies were determined positive.
- Through the theoretical calculations, we have (38) state with the total angular momentum and parity that have not symmetry by another practical value thus far.

3. 70Zn

Table 2. Excitation energy predictions for the 70Zn isotope using f5pvh interaction and observed experimental energies comparison

	tical values	Experimental values		
\mathbf{J}^{+}	E (MeV)	E (MeV)	J	
01	0.000	0.0	0+	
21	1.283	0.884	2+	
41	2.652	2.693	4+	
22	2.723	2.805		
11	3.043	2.949	1+	
02	3.339	3.328	(0+)	
23	3.356	3.634	2+	
31	3.411	3.411 3.419		
12	3.688	3.750	(0-, 1-, 2-)	

24	4.294	4.291	2+
25	4.576	4.588	(5, 6, 7-)
32	4.612	4.367	3+
03	4.674		
42	4.685	4.444	3+, 4+, 5+

By comparing the experimental data for this isotope in the table above with our theoretical findings using the f5pvh interaction, the following can be observed:

- A comparison with the available empirical data demonstrated that the ground state parity and total angular momentum of the 0+ level were identical.
- The angular momentum and parity of practical energy was determined to be (2.805,3.419)MeV, associated with angular momentum state 2^+ and 3^+,and(4.588) MeV with angular momentum 2 but negative parity. reflects the degree of alignment between the practical value and our theoretical value.
- By juxtaposing the theoretically calculated energies (1.283MeV; 2^+),(2.652MeV; 4^+),(3.043MeV; 1^+),(3.356MeV;2^+),(4.294MeV; 2^+),(4.612MeV; 3^+) and (4.685MeV;4^+) with the existing experimental data, we achieved a commendable concordance for the angular momenta.
- The parity and total angular momentum of the experimentally un confirmed energies (3.328) MeV, corresponding to positive angular momentum value of 2, And (3.750) MeV with angular momentum 2 but negative parity. This, are confirmed
- According to our calculations, the maximum experimental energy value is 6.116MeV, and the greater predicted energy is theoretically 4.685MeV.

3.2 Electromagnetic transition probability B(E2) and B(M1)

It is possible to think of gamma-rays as a type of electromagnetic radiation in which the electric field changes, causing the magnetic field to change as well. Radiation can be created by a fluctuating external magnetic field caused by an oscillating charge, or it can be produced by a changing magnetic field caused by a change in current or magnetic moment.

1. 70Se

We observed a reasonable agreement between the available experimental data and the theoretical data electrical the transitions $B(E2)2_1^+ \rightarrow 0_1^+ , B(E) 2_2^+ \rightarrow 0_1^+, B(E)$ $2 \ 2^+ \rightarrow 2 \ 1^+, \ B(E) \ 4 \ 1^+ \rightarrow 2 \ 1^+ \ and \ B(E)$ $4 \quad 1^+ \rightarrow 2 \quad 2^+$ for using the (f5pvh) interaction. We observed a reasonable agreement between the available experimental data and the theoretical data for the electrical transitions B(M1) 2 $2^+ \rightarrow 2 1^+$ for using the (f5pvh) interaction. Our calculations revealed new transitions that had no experimental values before. A satisfactory agreement was observed for the transitions B(M1) 2 $2^+ \rightarrow 2 1^+$ when we compared some of our findings for the magnetic transitions of the interaction (f5pvh) with the experimental study (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Data Sheets 136, 1 (2016)10.1016/j.nds.2016.08.001" n.d.).

Table 4. Comparison of the B(E2) and B(M1) results by using) f5pvh (interaction in units e2fm4 and μ _N^2 respectively for the 70Se isotope with the experimental data

7+	\rightarrow	L	$B(M1) \left(\mu_N^2\right)$		$B(E2) e^2 fm^4$	
J_i^+		J_f	Theory	Exp.	Theory	Exp.
21	\rightarrow	01+	0.0000		316.500 0	17.9
02	\rightarrow	21+	0.0000		85.6000	
22	\rightarrow	01+	0.0000		4.3210	0.19
22	\rightarrow	2+	0.0024	0.009	266.300 0	33
22	\rightarrow	0+	0.0000		142.000 0	
23	\rightarrow	01+	0.0000		8.3950	
23	\rightarrow	21+	0.0366		27.7100	
23	\rightarrow	02+	0.0000		161.800 0	
23	\rightarrow	2+	0.1371		3.3420	
41	\rightarrow	2+	0.0000		438.200 0	21.5
41	\rightarrow	2+	0.0000		7.4550	17
4 ₁	\rightarrow	2*	0.0000		9.9710	

2. 70Zn

We observed a reasonable agreement between the available experimental data and the theoretical data for the electrical transitions B(E2) $2_1^+ \rightarrow 0_1^+$, B(E2) $4_1^+ \rightarrow 2_1^+$, B(E2) $2_2^+ \rightarrow 0_1^+$, B(E2) $2_2^+ \rightarrow 2_1^+$, B(E2) $1_1^+ \rightarrow 2_2^+$ and B(E2)

 $0_2^+ \rightarrow 2_1^+$ for using the (f5pvh) interaction. Our calculations revealed new transitions that had no experimental values before.

A satisfactory agreement was observed for the transitions B(M1) $2_2^+\rightarrow 2_1^+$ and B(M1) $1_1^+\rightarrow 2_2^+$ when we compared some of our findings for the magnetic transitions of the interaction (f5pvh) with the experimental study (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Data Sheets 136, 1 (2016) 10.1016/j.nds.2016.08.001" n.d.)

Table 5. Comparison of the B(E2) and B(M1) results by using) f5pvh (interaction in units e2fm4 and μ_N^2 respectively for the 70Zn isotope with the experimental data

7+	\rightarrow	1.	$B(M1) \left(\mu_N^2\right)$		B(E2)	$e^2 f m^4$
J_i^+	\rightarrow	J_f	Theory	Exp.	Theory	Exp.
21	\rightarrow	0+	0.0000		114.600 0	286.1763
41	\rightarrow	21+	0.0000		79.4500	325.5898
22	\rightarrow	01+	0.0000		9.8500	10.2818
22	\rightarrow	2+	0.4488	0.017	38.4200	171.3630
22	\rightarrow	41+	0.0000		72.8900	
11	\rightarrow	01+	0.0872		0.0000	
11	\rightarrow	21+	0.0036		2.1860	
11	\rightarrow	2+	3.5930	0.003 9	2.9340	188.4993
02	\rightarrow	2+	0.0000		0.1468	639.1841

3.70Ga

We observed a reasonable agreement between the available experimental data and the theoretical data for the electrical transitions B(E2) 3 $1^+ \rightarrow 2 1^+$ for using the (f5pvh) interaction. Our calculations revealed new transitions that had no experimental values before. A satisfactory agreement was observed for the transitions B(M1) 3 $1^+ \rightarrow 2 1^+$ when we compared some of our findings for the magnetic transitions of the interaction (f5pvh) with the experimental study (No Title13. "1-Jul-2016 ENSDF insertion: 2016-09 Publication: Nuclear Sheets 136, 1 (2016)10.1016/j.nds.2016.08.001" n.d.).

Table 2. Comparison of the B(E2)and B(M1) results by using) f5pvh (interaction in units e2fm4 and μ _N^2 respectively for the 70Ga isotope with the experimental data

I+	L	B(M1)	$\left(\mu_N^2\right)$	B(E2)	$e^2 fm^4$	
Ji) ;	Theory	Exp.	Theory	Exp.

-	\rightarrow	11	0.0166		25.860	
-	\rightarrow	21	0.0000		33.7600	
-	\rightarrow	11	0.0000		2.0350	
	\rightarrow	21	0.4340	1.253 0	0.0007	0.0027
_	\rightarrow	41	0.8921		19.0200	
-	\rightarrow	11	0.0000		84.6300	
-	\rightarrow	21	1.1970		15.7200	
-	\rightarrow	41	0.1289		4.5750	
-	\rightarrow	31	0.3053		39.0800	

4. Conclusions

This study conducted extensive shell model computations for neutron-rich, isotopes of (Se, Ga, Zn) for A=70. The current work demonstrated that the energy levels and electromagnetic transition probability were computed using the interaction f5pvh and the results were shown to be in reasonable accord with the existing experimental data. Numerous energy stats have been confirmed using the interaction and additional energy levels have been acquired in our computations. The B(E2) and B(M1) values also showed some degree of compatibility with the experimental results.

Author Statements:

- **Ethical approval:** The conducted research is not related to either human or animal use.
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